



Plasma Sheath Around Long Conductors with Elliptic Cross- sections

L. Chiabò, G. Sánchez-Arriaga

Universidad Carlos III de Madrid, Spain

APS-DPP 2020, 9-13 November

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- Introduction
- Kinetic Model for 2D Emissive Probes with Elliptic cross-section
- Numerical Results
- Conclusions

INTRODUCTION

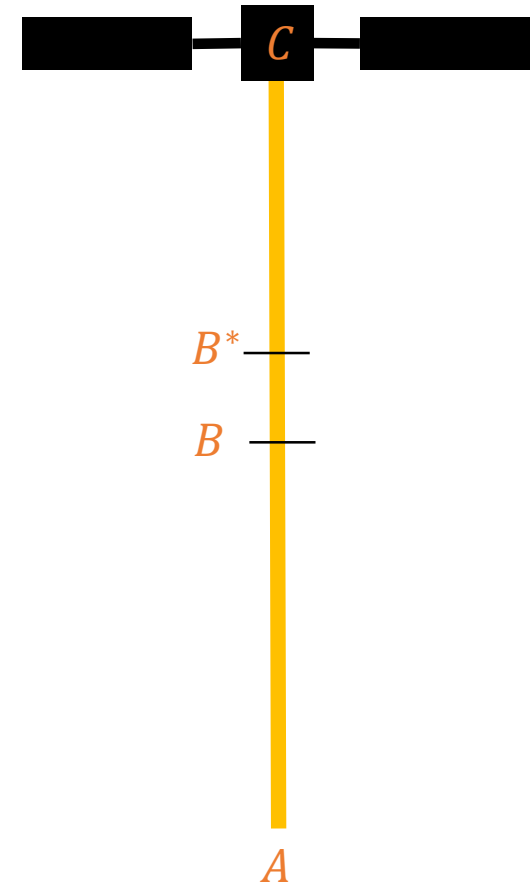
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- Goal: Kinetic analysis of plasma structure and currents around long electron emitting objects with elliptical cross-sections

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- Applications
 - Low Work-function Tethers (E.T.PACK)

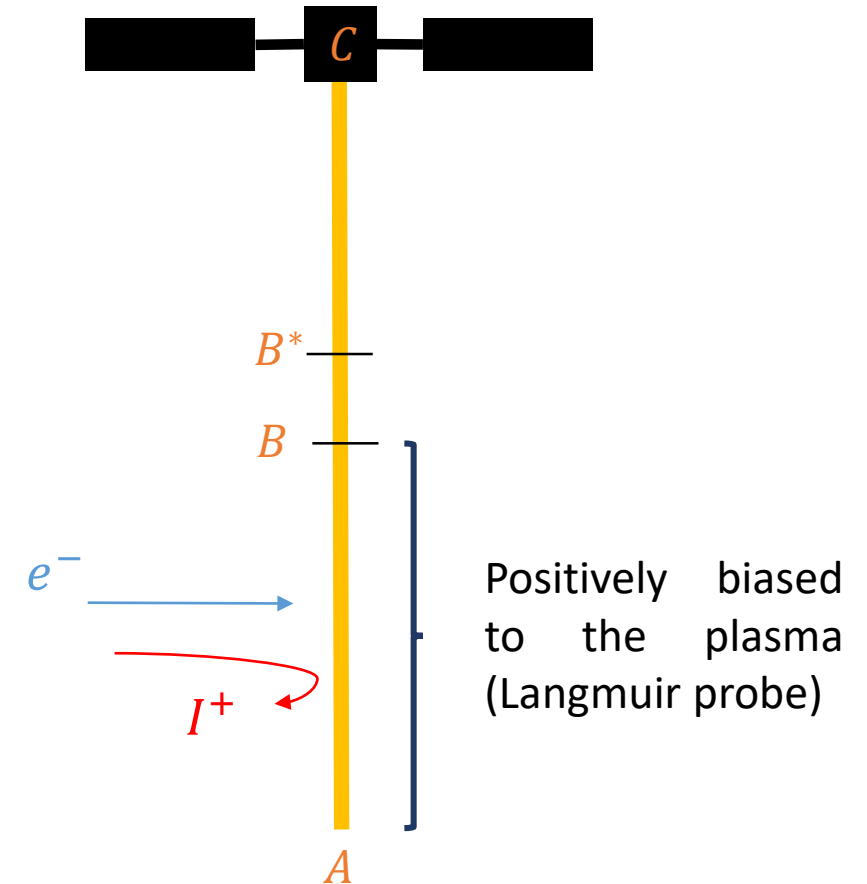
Plasma-Tether interaction



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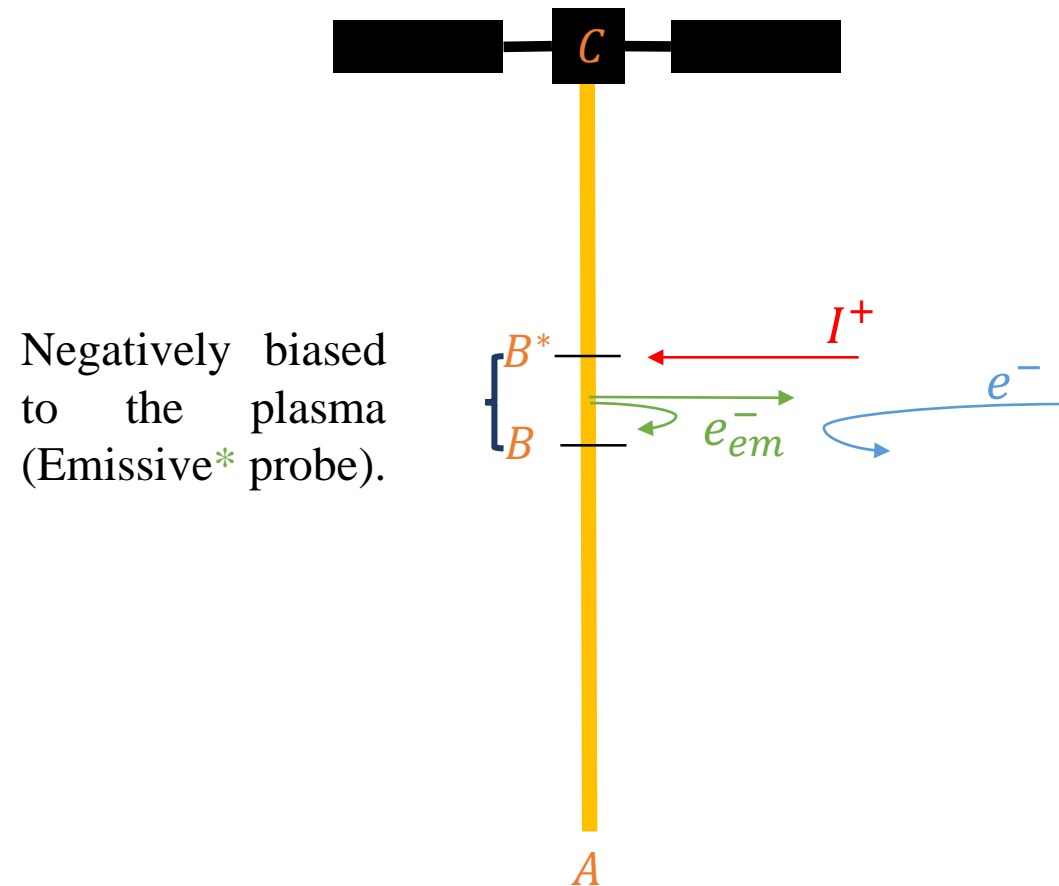
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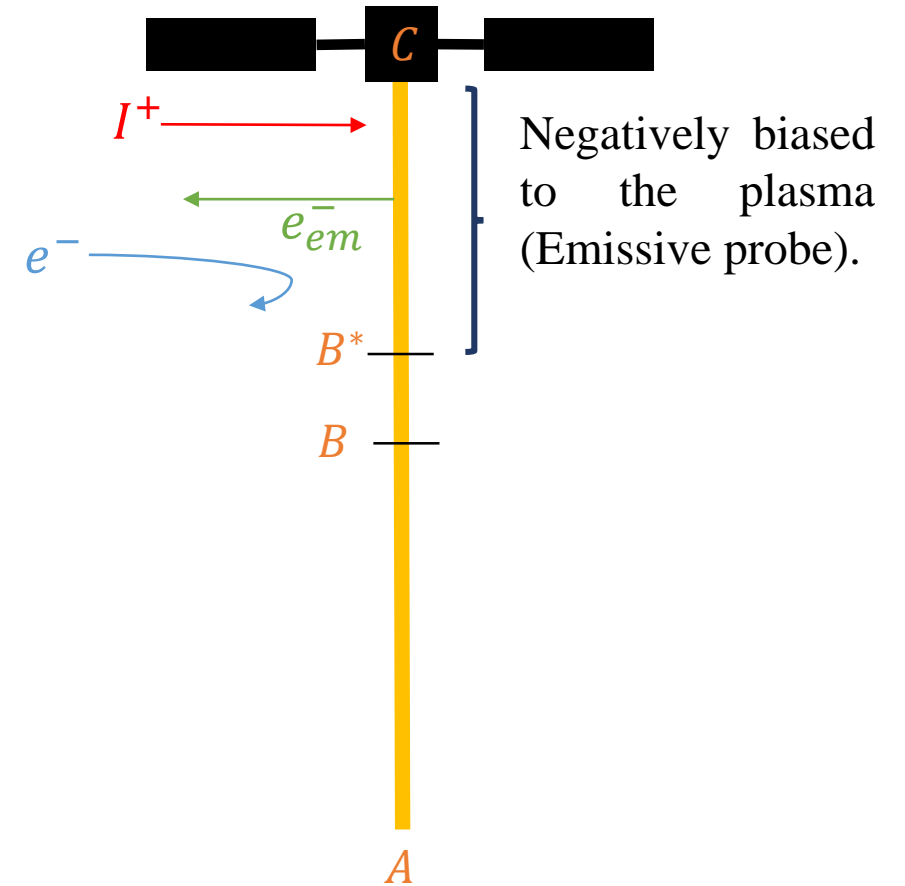


*Space Charge Limited (SCL): a potential dip reflects part of the emitted electrons back to the tether

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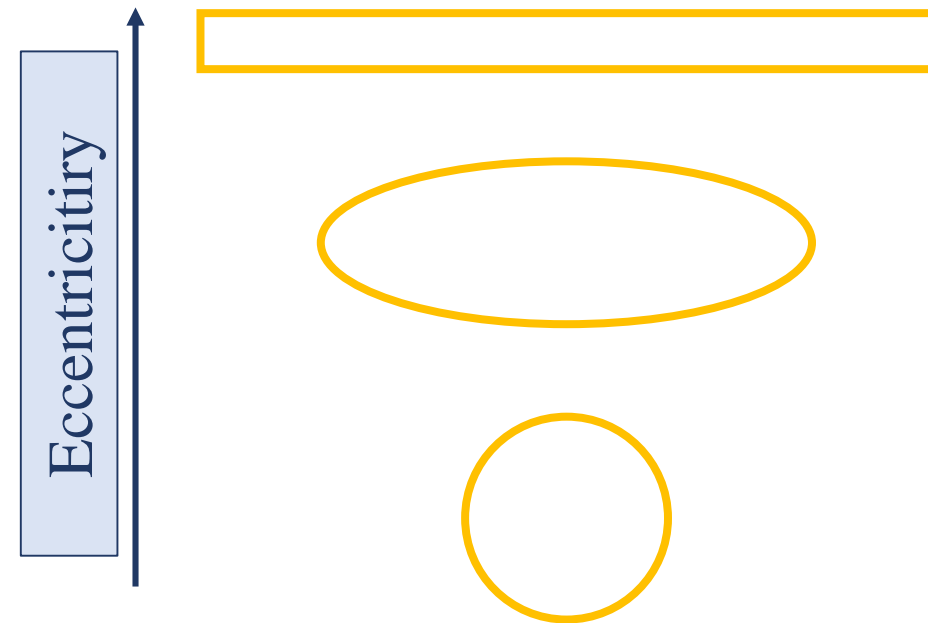


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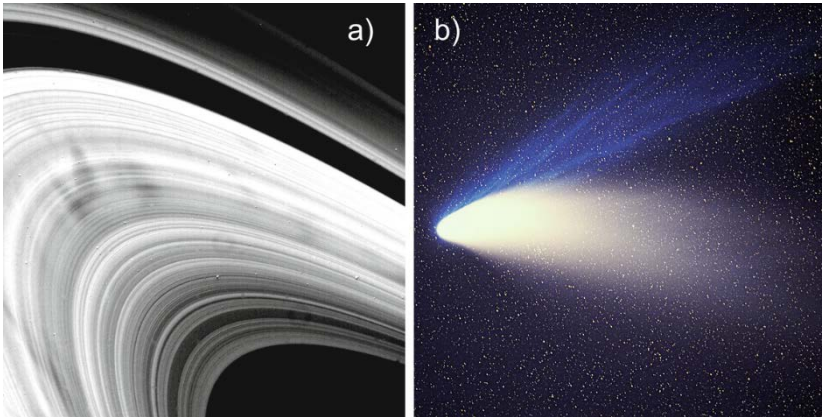
Plasma-Tether interaction

- Tape geometry: favourable for tethers applications
- Ellipse (eccentricity = 1) \approx Tape



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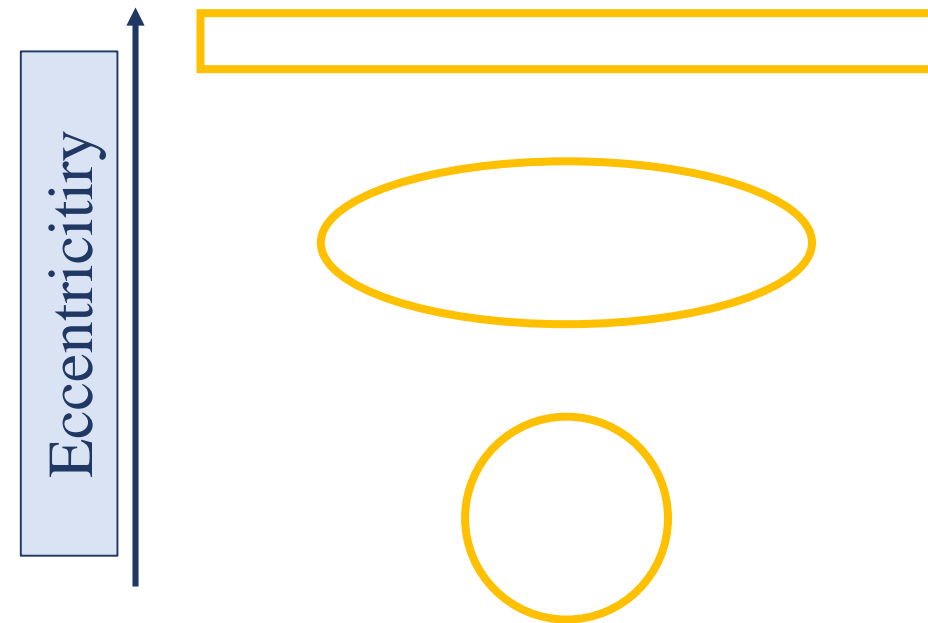
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 - Low Work-function Tethers (E.T.PACK)
 - Probe theory
 - Dusty plasmas
 - Spacecraft charging



Melzer A. (2019) Introduction. In: *Physics of Dusty Plasmas. Lecture Notes in Physics*, vol 962. Springer, Cham.

Plasma-Tether interaction

- Tape geometry: favourable for tethers applications
- Ellipse (eccentricity = 1) \approx Tape



KINETIC MODEL FOR 2D
EMISSIVE PROBES WITH
ELLIPTIC CROSS-SECTION

KINETIC MODEL FOR 2D EMISSIVE PROBES

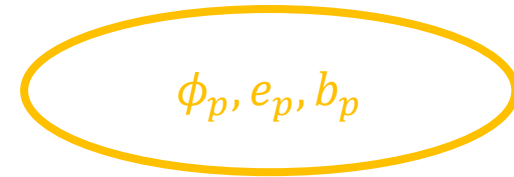
Hypotheses

- Stationary solution

KINETIC MODEL FOR 2D EMISSIVE PROBES

Hypotheses

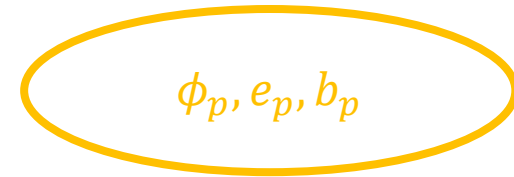
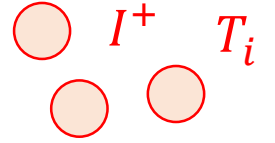
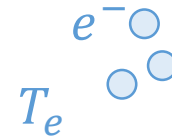
- Stationary solution
- Probe at rest in an unmagnetized, collisionless, Maxwellian plasma with background density N_0


$$\phi_p, e_p, b_p$$

KINETIC MODEL FOR 2D EMISSIVE PROBES

Hypotheses

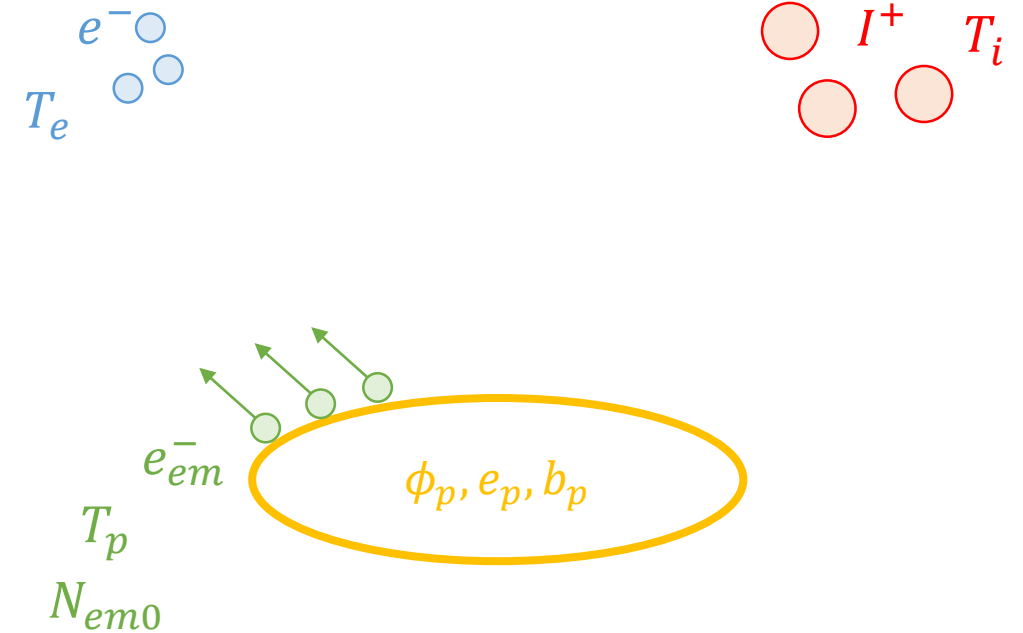
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KINETIC MODEL FOR 2D EMISSIVE PROBES

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- Half-Maxwellian e_{em}^-
- Emission follows Richardson-Dushman law



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- Stationary solution
- Probe at rest in an unmagnetized, collisionless, Maxwellian plasma with background density N_0
- Singly charged I^+
- Half-Maxwellian e_{em}^-
- Emission follows Richardson-Dushman law
- No trapped particles
- Our solutions depend on 6 parameters:

e_p : probe eccentricity

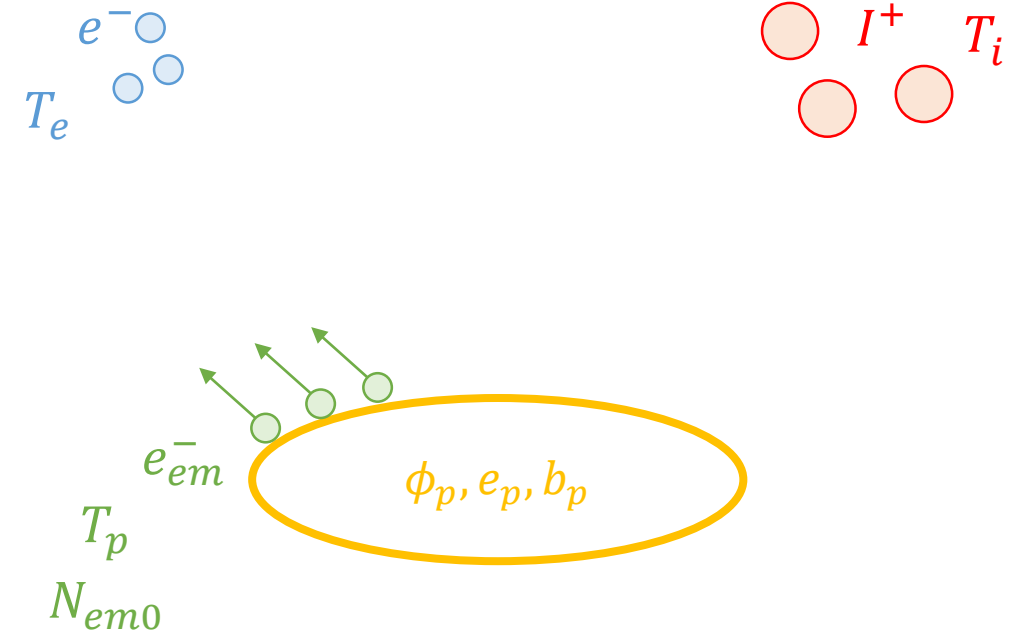
b_p : semi minor axis/Debye length

ϕ_p : normalized probe bias

$$\delta_p \equiv T_p/T_e$$

$$\delta_i \equiv T_i/T_e$$

$$\beta \equiv \frac{N_{em0}}{N_0} = \text{emission level}$$



Vlasov-Poisson system

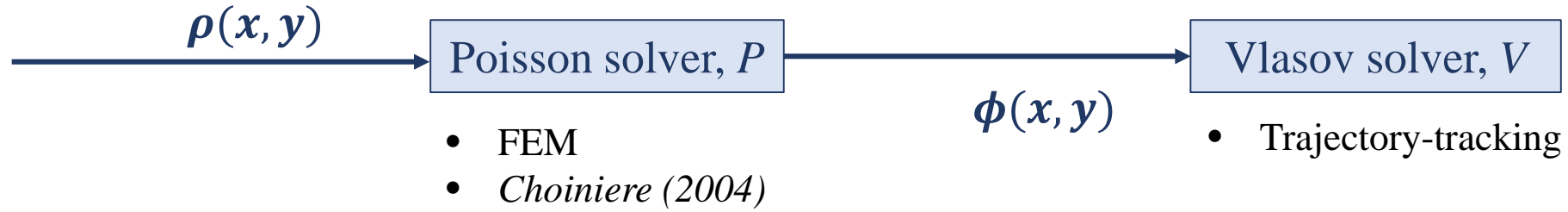
NUMERICAL SCHEME

$\rho(x, y)$

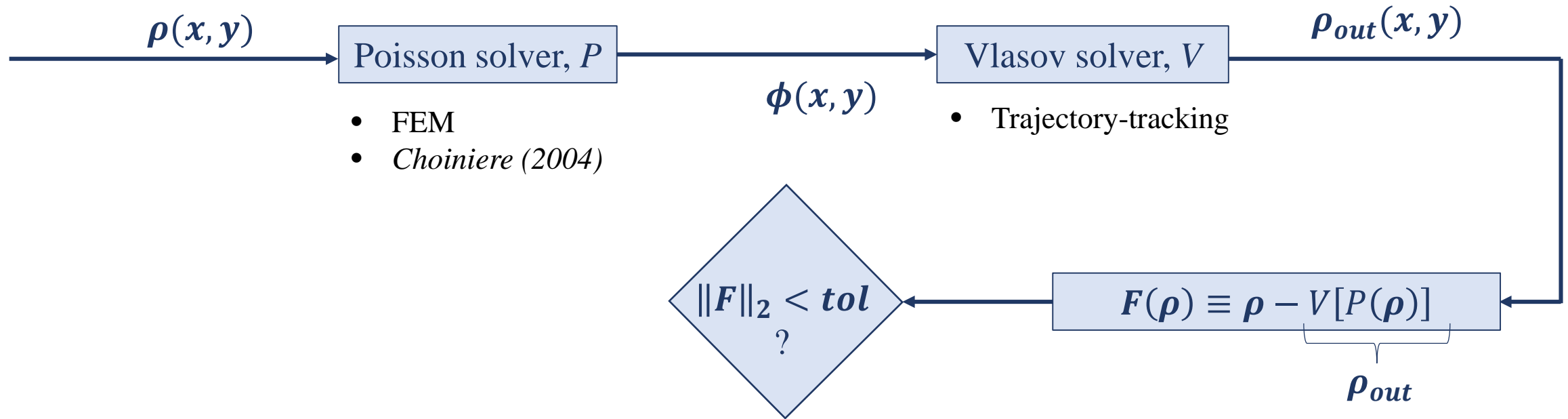
→ Poisson solver, P

- FEM
- *Choiniere (2004)*

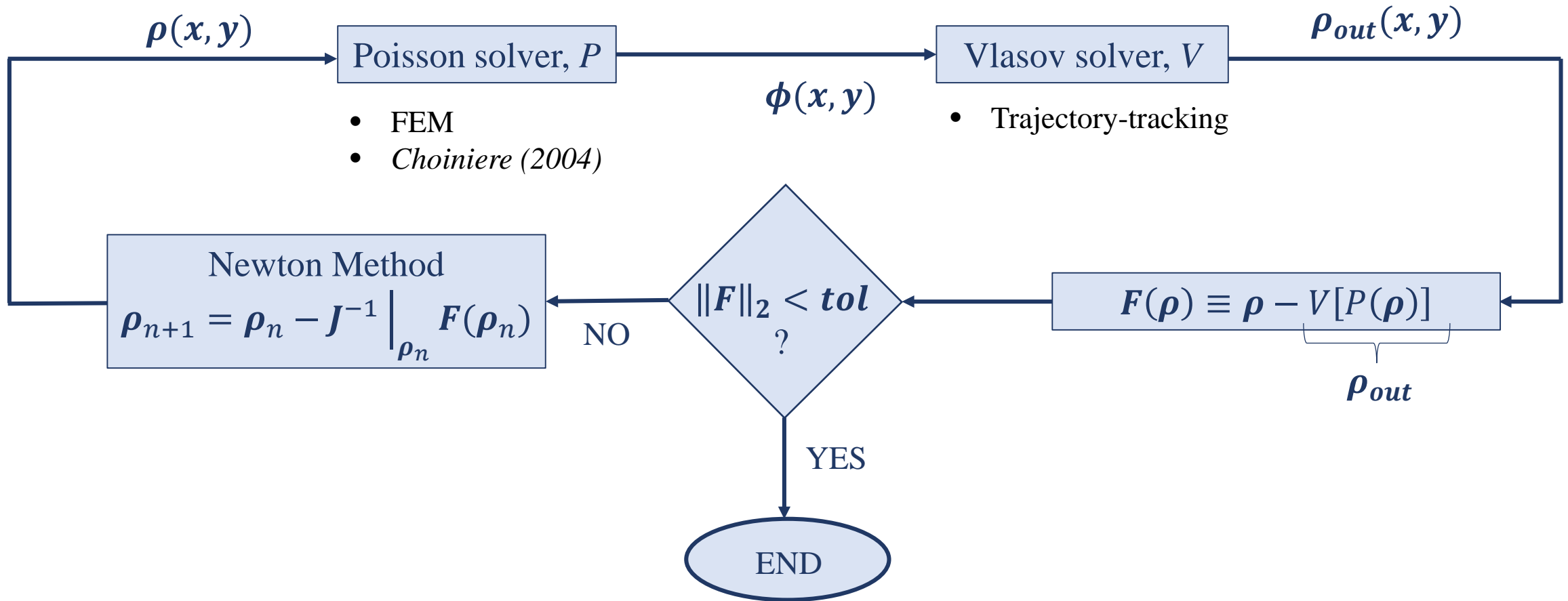
NUMERICAL SCHEME



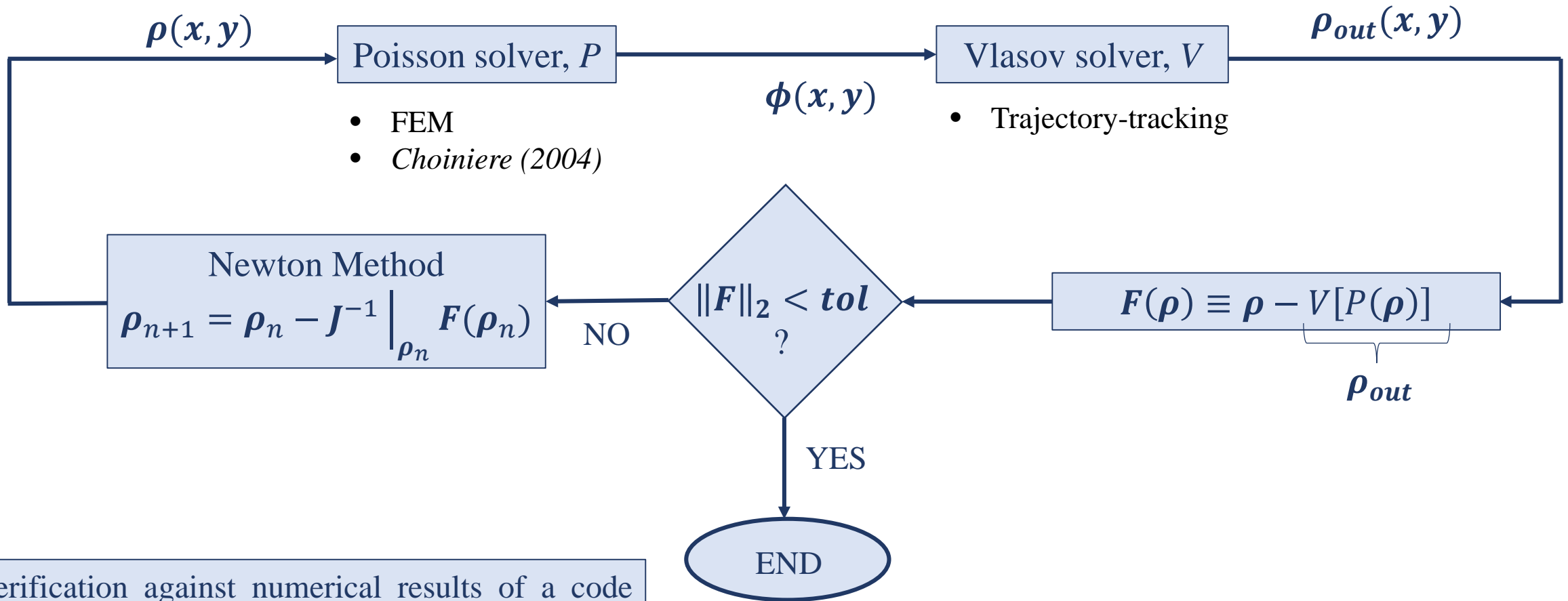
NUMERICAL SCHEME



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Verification against numerical results of a code for cylindrical emissive probes by *Chen, Sánchez-Arriaga (2017)*

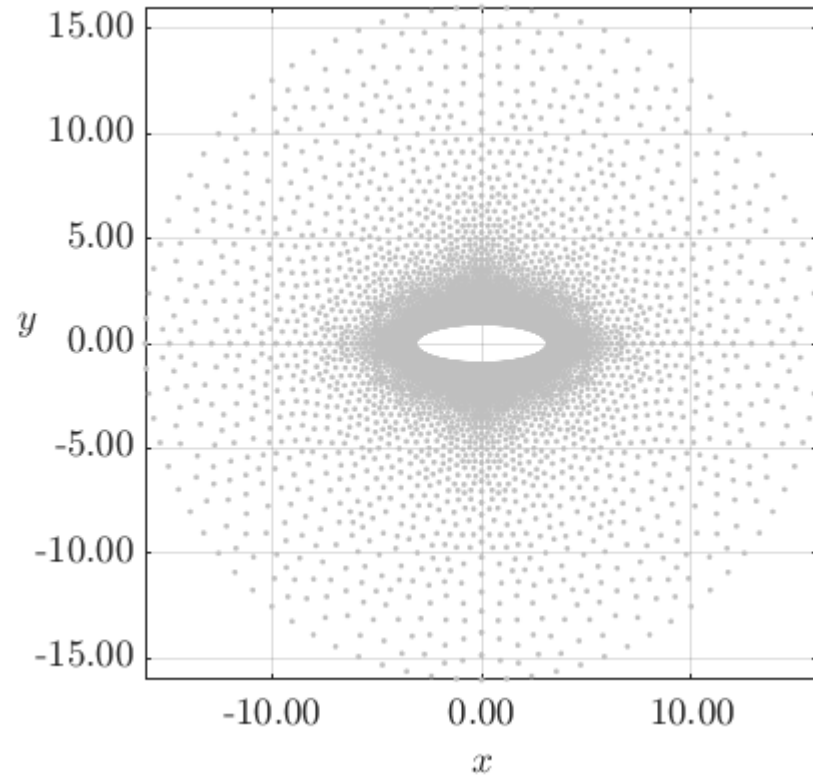
NUMERICAL SCHEME

Vlasov solver

- Computes the density of each species (ϕ as input)
- Collisionless plasma: $df_{\alpha}/d\tau = 0$
- Backward integration of particles orbits

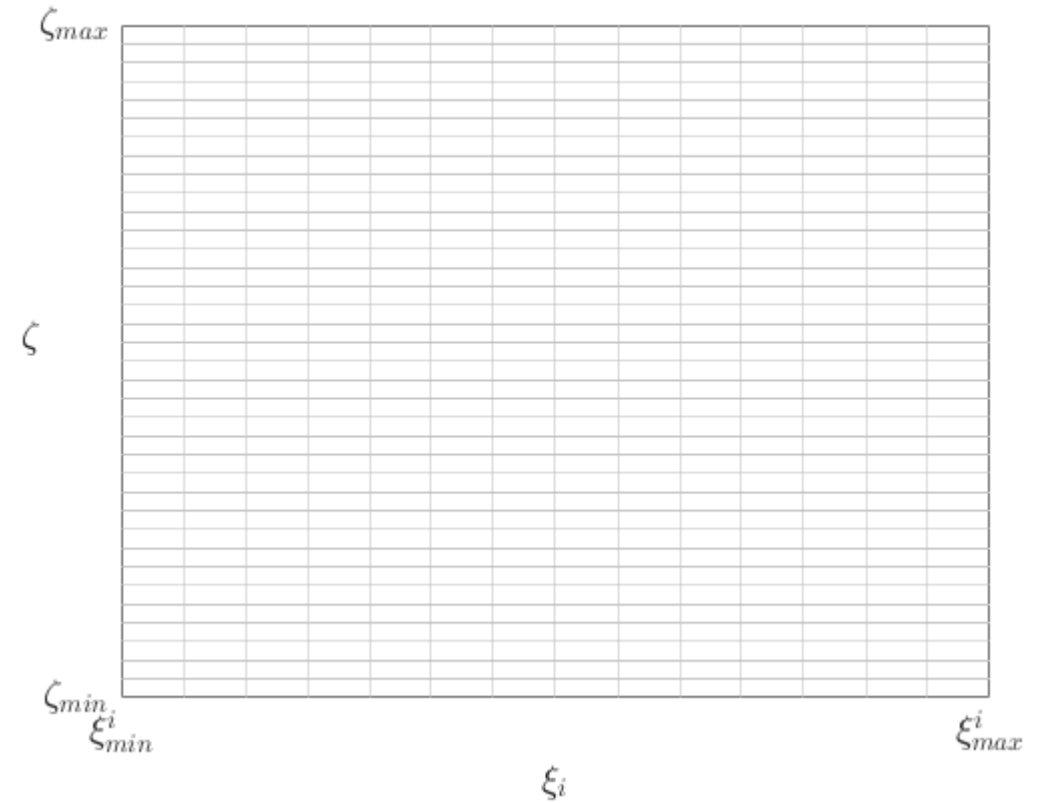
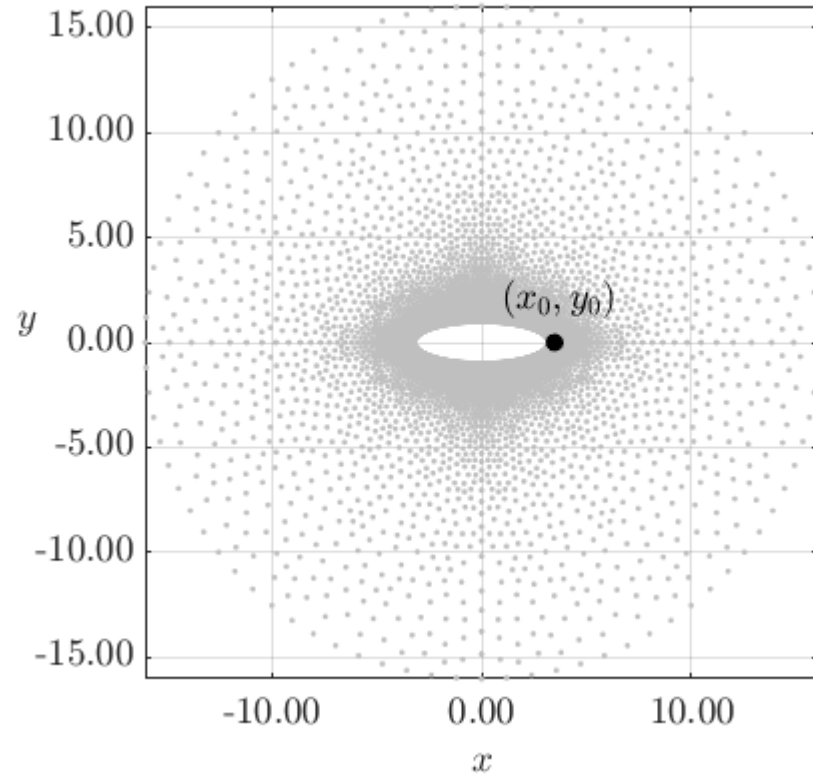
NUMERICAL SCHEME

Trajectory-tracking algorithm



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Trajectory-tracking algorithm



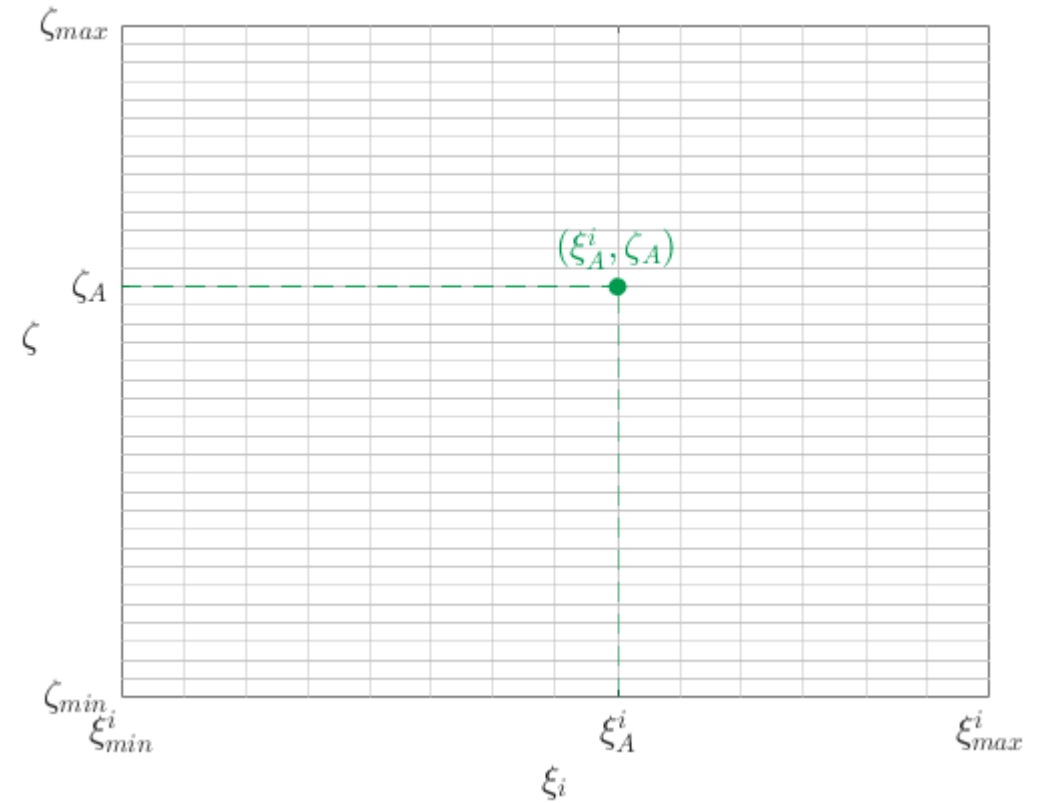
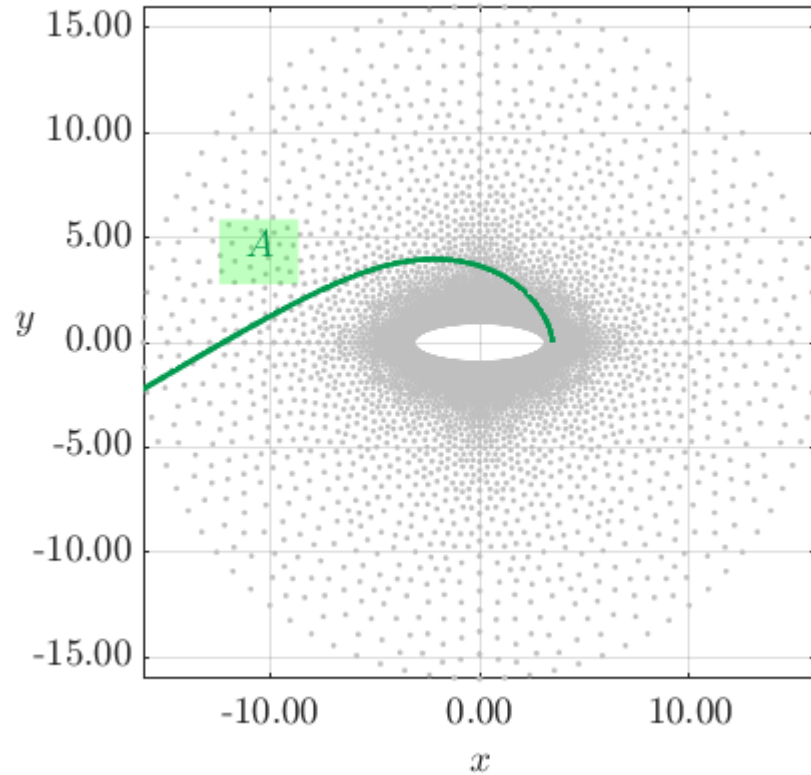
Plasma ions around a negatively polarized probe

$\xi_i \equiv$ dimensionless energy, $\zeta \equiv$ velocity angle

NUMERICAL SCHEME

Trajectory-tracking algorithm

- Finite value (Maxwellian)
- Null value



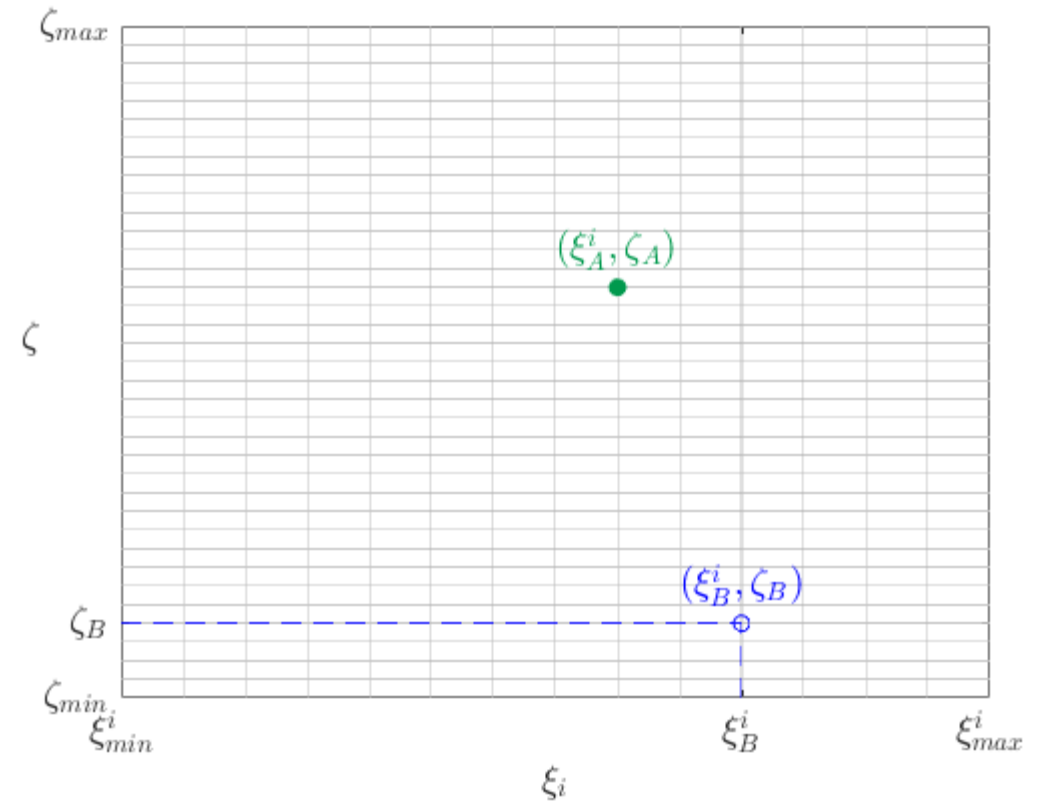
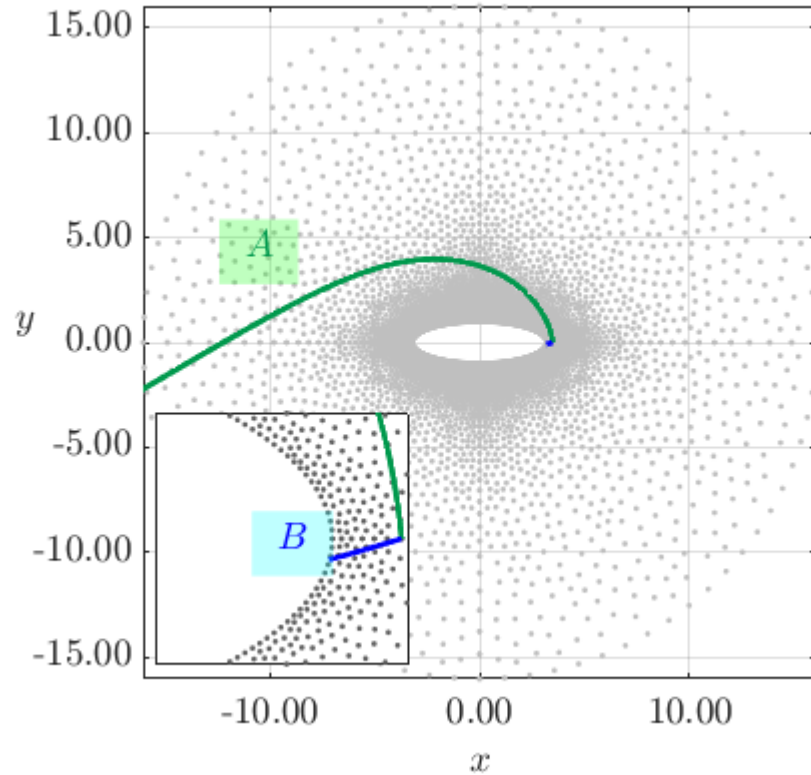
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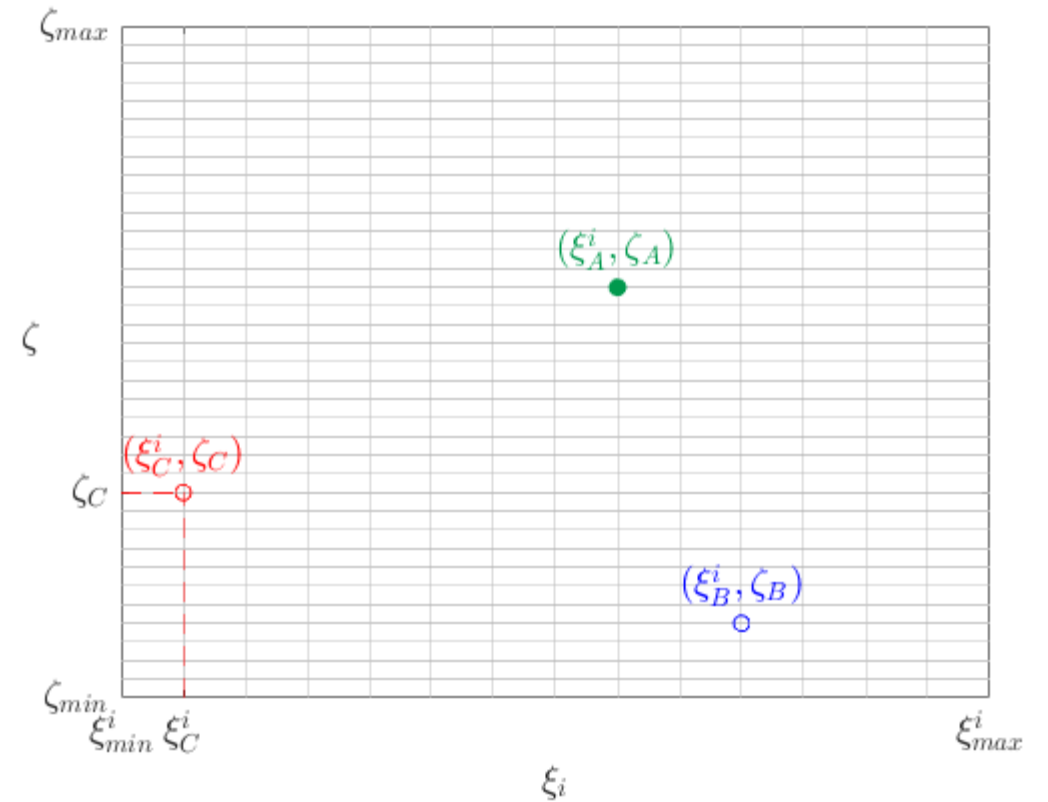
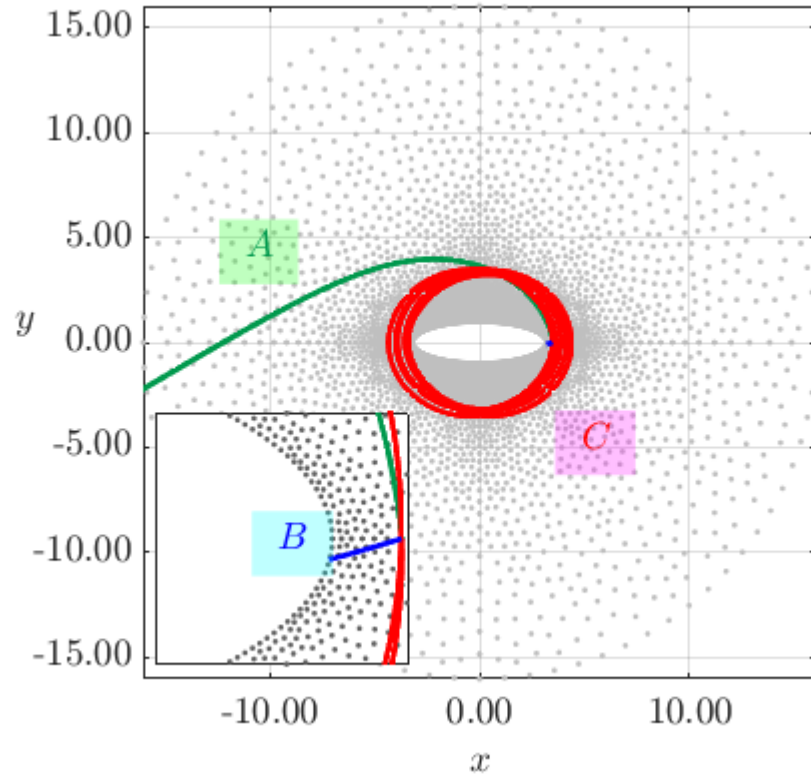
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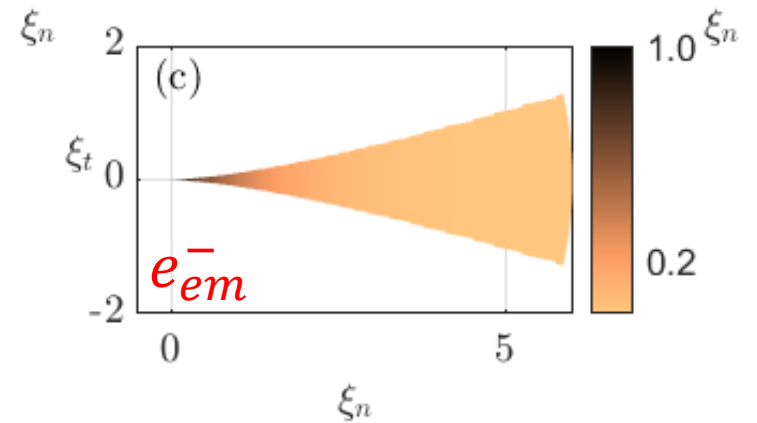
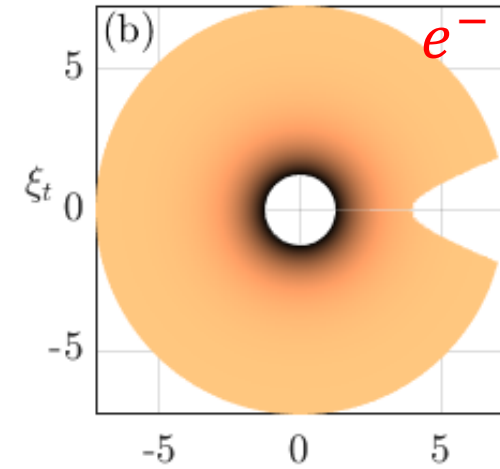
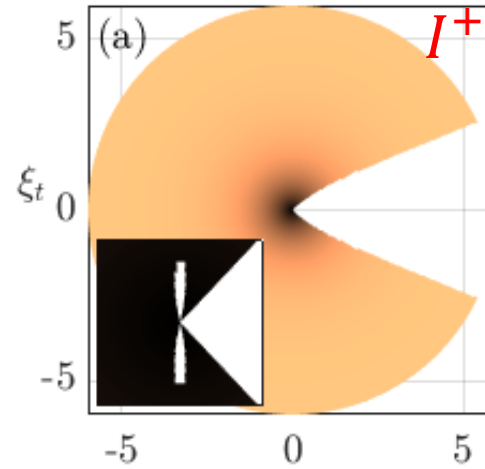
Vlasov solver

- Computes the density of each species (ϕ as input)
- Collisionless plasma: $df_\alpha/d\tau = 0$
- Backward integration of particles orbits
- Numerical integration:

$$n_\alpha(x_0, y_0) = \sum_{i=1}^{N_\xi} \sum_{j=1}^{N_\zeta} f_\alpha(x_0, y_0, \xi_i^\alpha, \zeta_j) \Delta\zeta_j \Delta\xi_i^\alpha$$

$$\xi_n = \xi_\alpha \cdot \cos \zeta$$

$$\xi_t = \xi_\alpha \cdot \sin \zeta$$



NUMERICAL RESULTS

LIMITATIONS OF STATIONARY VP-SOLVERS

“Limitations of Stationary Vlasov-Poisson Solvers in Probe Theory” (under review)

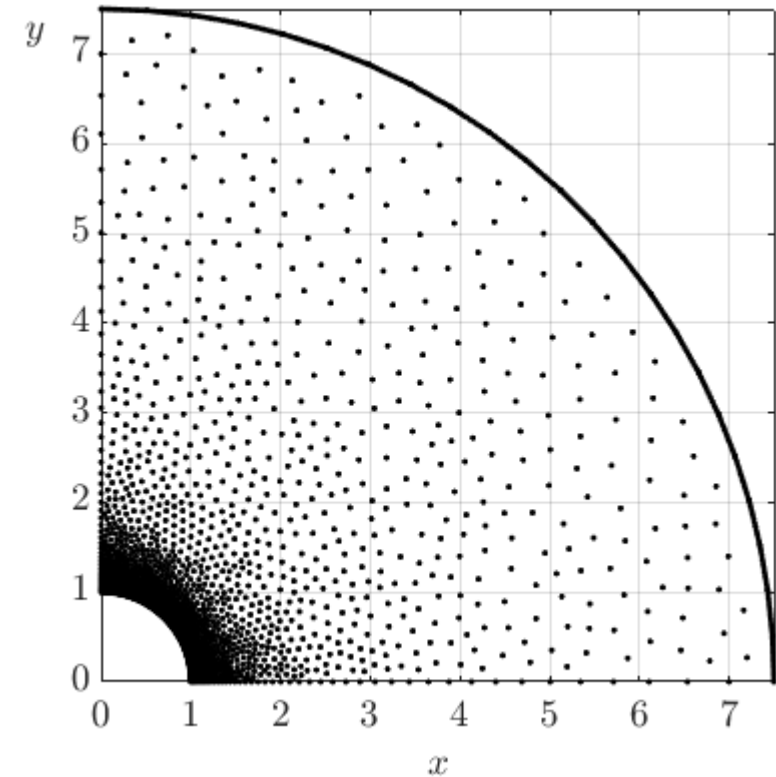
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“*Limitations of Stationary Vlasov-Poisson Solvers in Probe Theory*” (under review)

$$\phi_p = -4, e_p = 0, b_p = 1, \delta_i = 1$$

Numerical

- Angular momentum and energy are physical invariants for cylindrical probes



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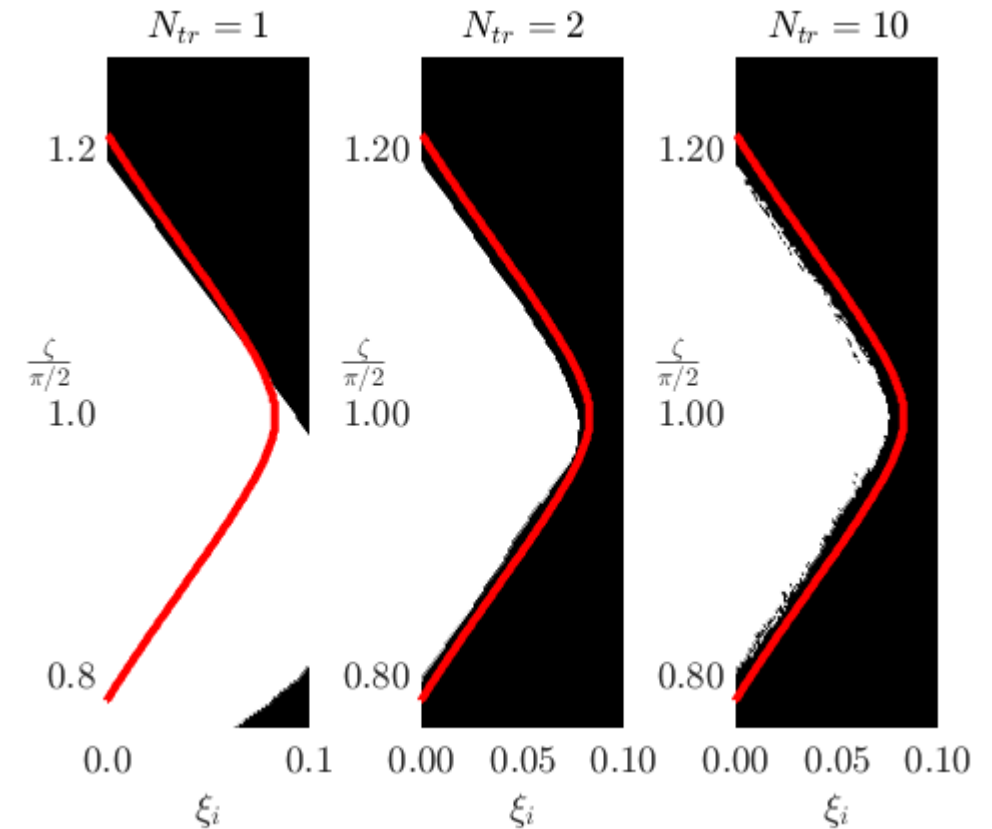
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N_{tr} \equiv number of complete turns a around the probe for trapped particles



Red boundary: computed analytically with code by *Chen, Sánchez-Arriaga(2017)*

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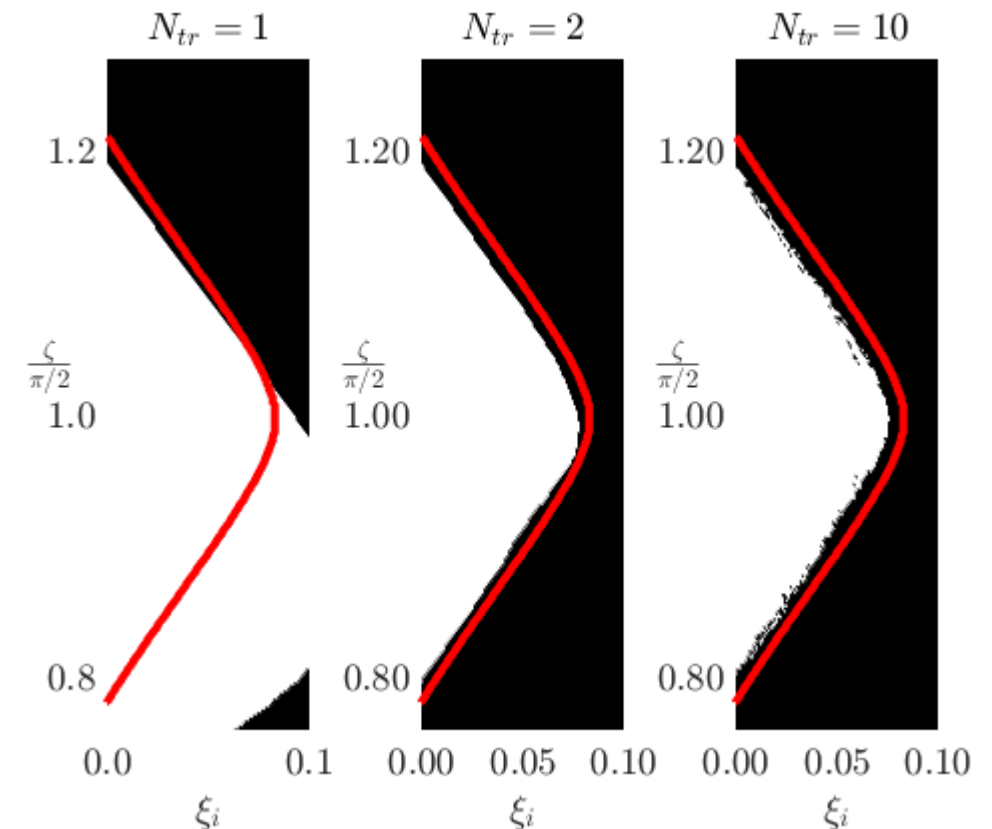
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Numerical

- Angular momentum and energy are physical invariants for cylindrical probes
- Spatial discretization of the electric field: angular momentum & total energy are not perfectly conserved by the algorithm
- Boundary between trapped (white) and non-trapped (black) particles becomes jagged as N_{tr} increases

$N_{tr} \equiv$ number of complete turns a around the probe for trapped particles



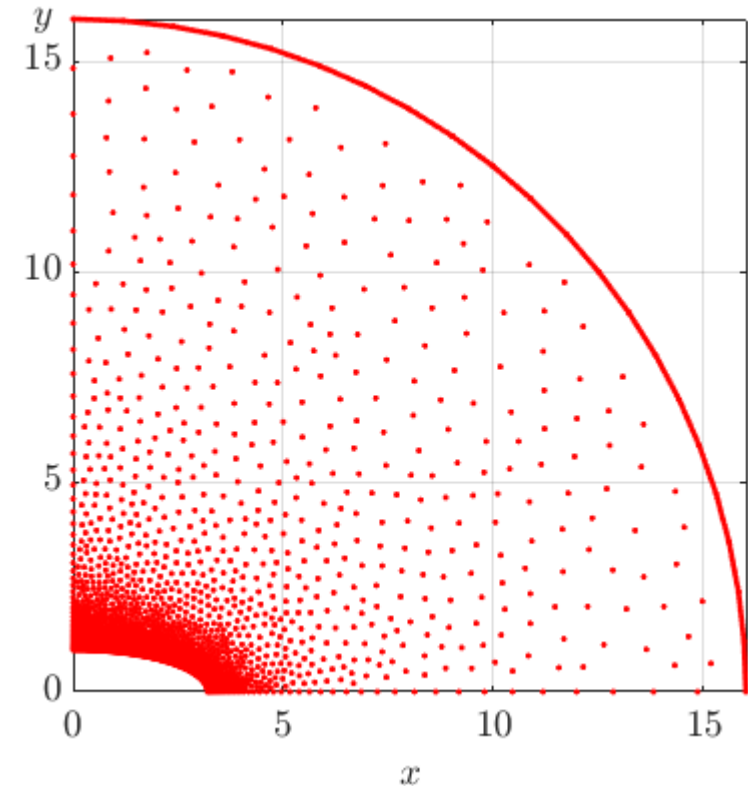
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$$\phi_p = -4, e_p = 0.95, b_p = 1, \delta_i = 1$$

Physical limitations for elliptic probes



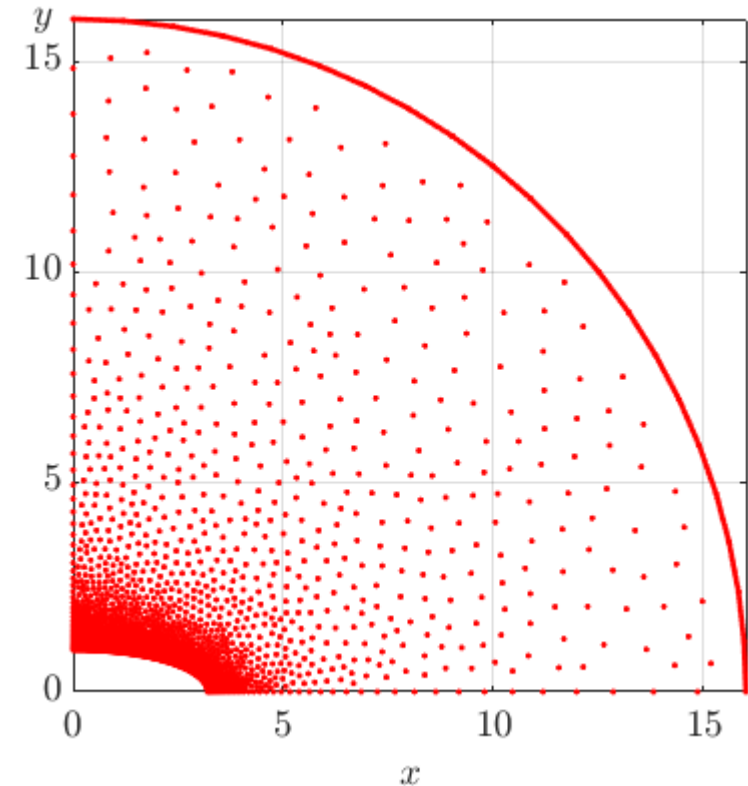
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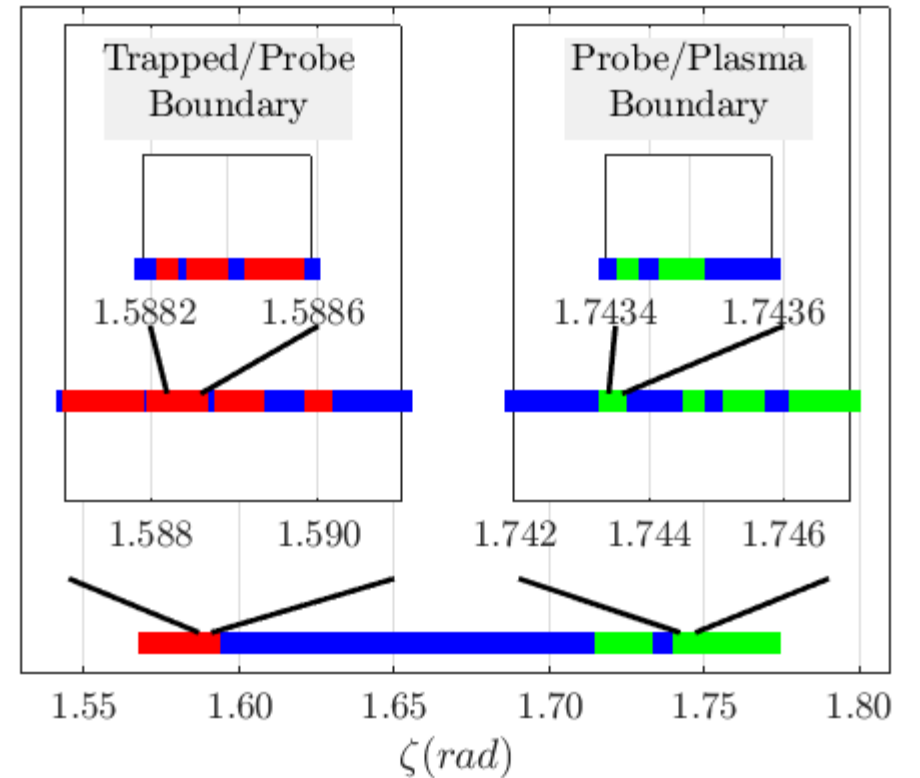
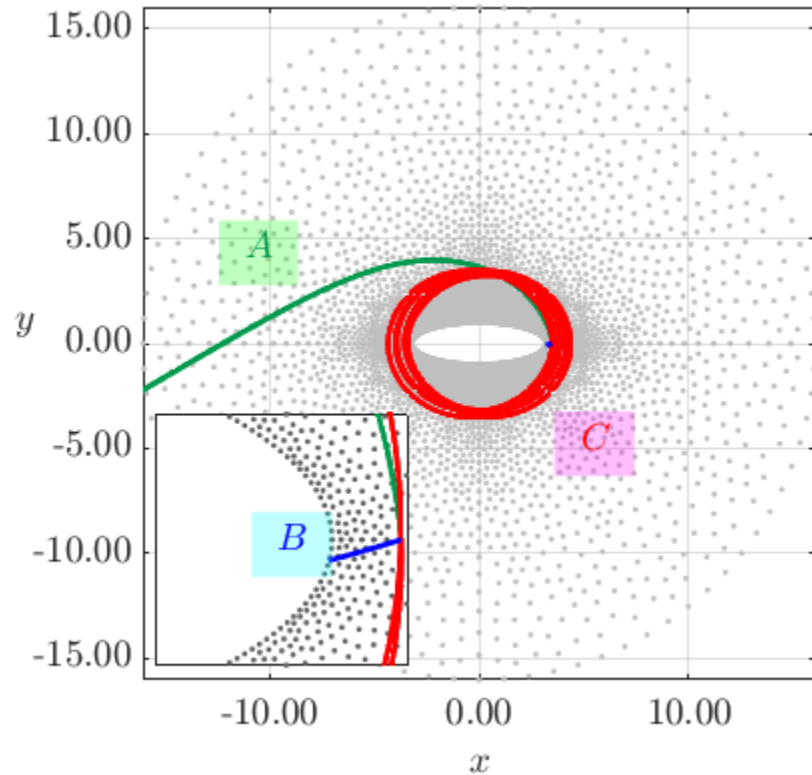
Physical limitations for elliptic probes

- Angular momentum is not conserved
- Chaotic orbits appear



LIMITATIONS OF STATIONARY VP-SOLVERS

“Limitations of Stationary Vlasov-Poisson Solvers in Probe Theory” (under review)



$$\xi_i \approx 0.062, \zeta \text{ varies}$$

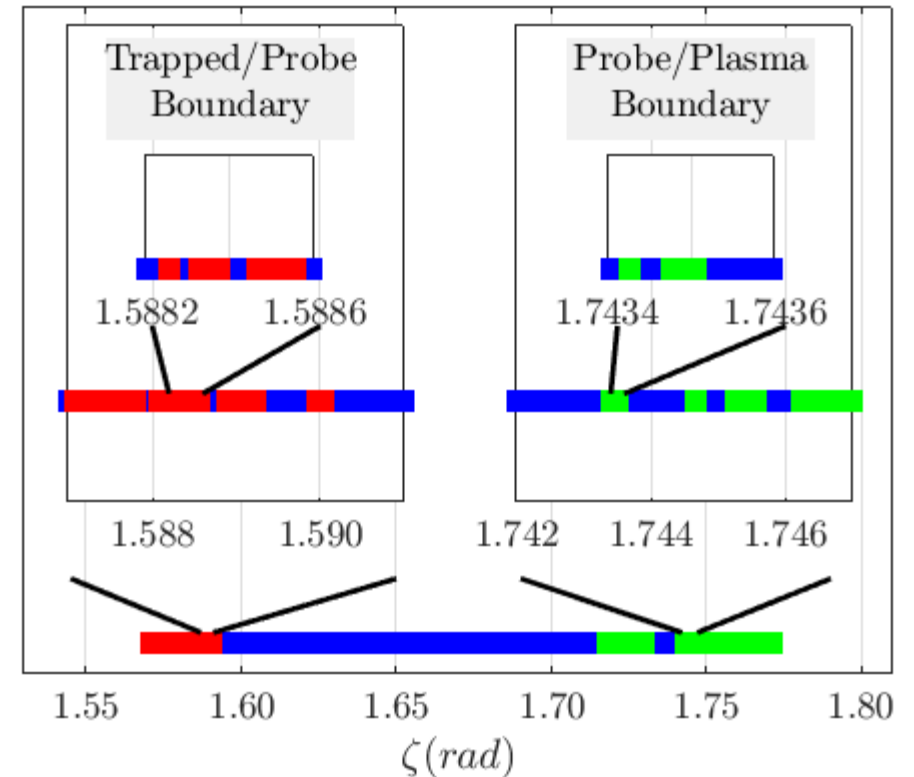
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Physical limitations for elliptic probes

- Angular momentum is not conserved
- Chaotic orbits appear
- Boundaries in velocity space reveal the distribution function is filamented



$$\xi_i \approx 0.062, \zeta \text{ varies}$$

SPACE CHARGE LIMITED (SCL) TRANSITION

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$$\phi_p = -4, e_p = 0.75, b_p = 1, \delta_i = 1, \delta_p = 0.32, \beta \in [0, 30]$$

- Impact of the emission level

$$\beta \propto \exp\left(-\frac{eW}{k_B T_P}\right)$$

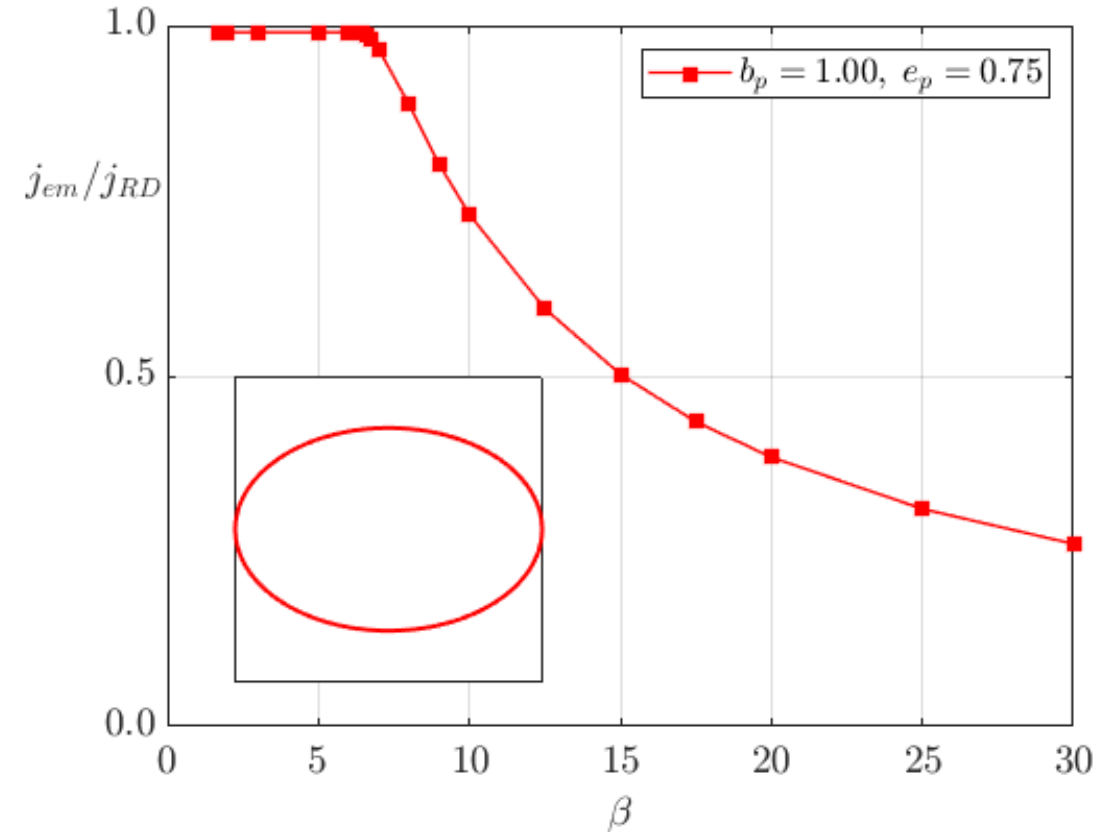
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- $\frac{j_{em}}{j_{RD}} = \frac{\text{Emitted electrons reaching infinity}}{\text{RD maximum}}$



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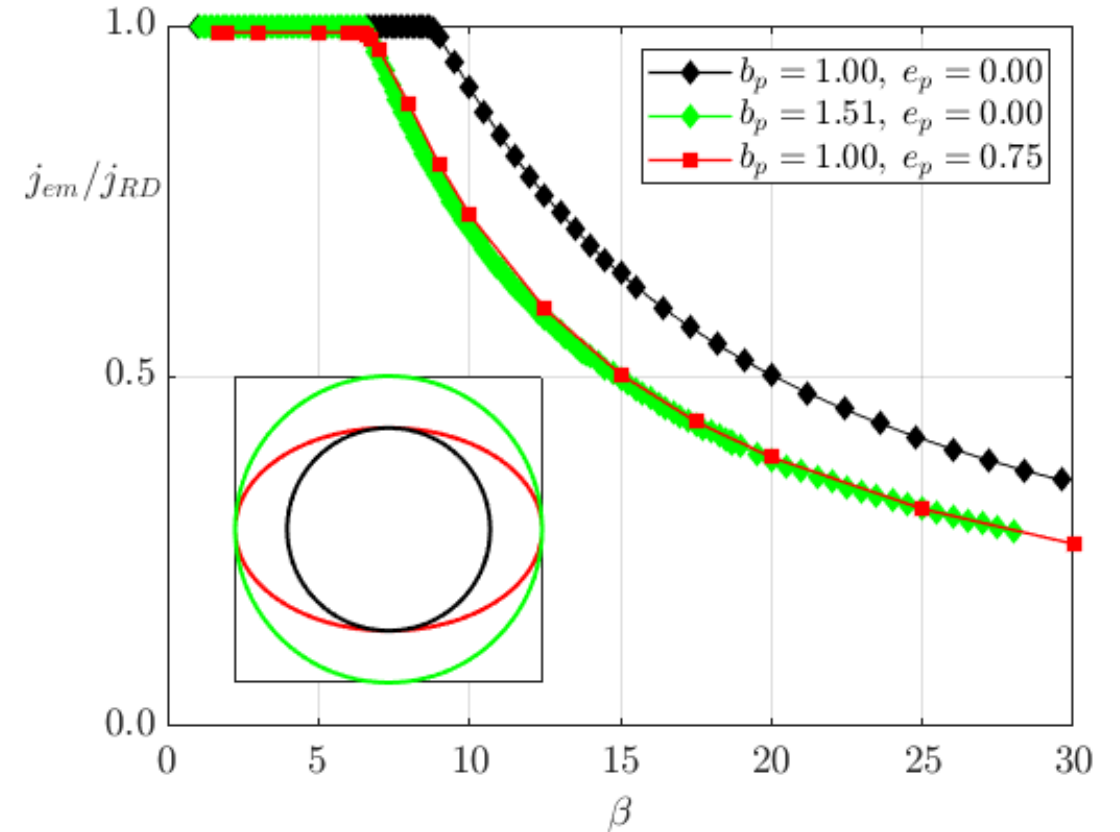
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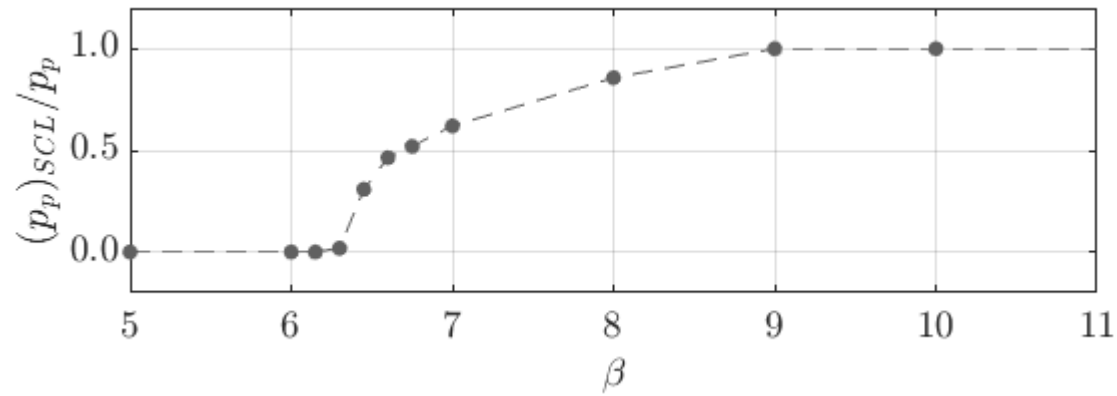
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Elliptical probe \approx Round probe with radius equal to the semi-major axis



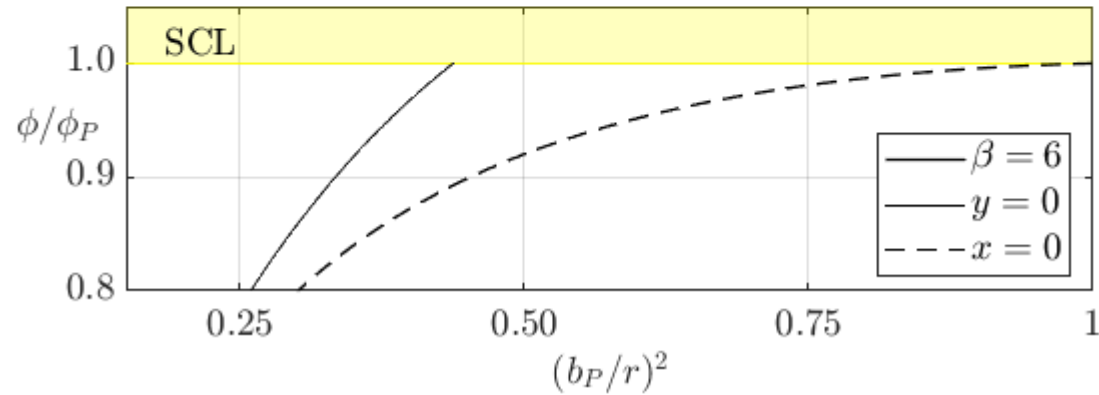
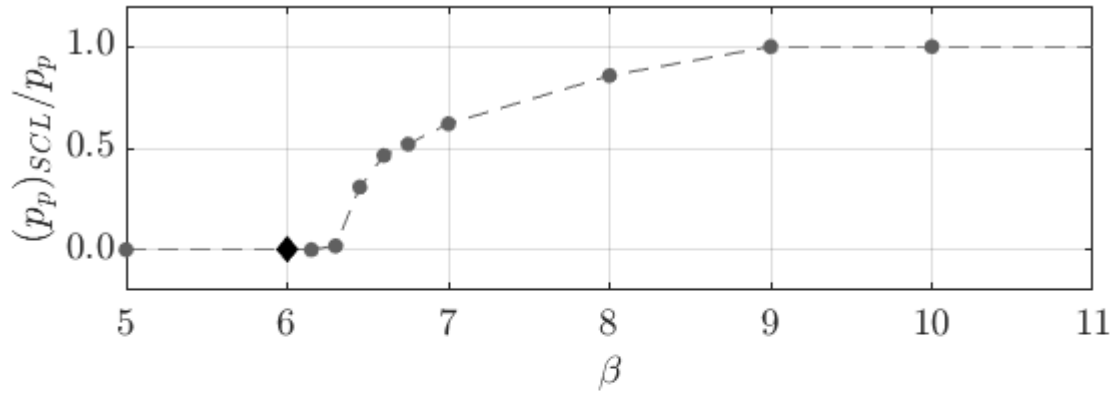
SPACE CHARGE LIMITED (SCL) TRANSITION

$$\frac{(p_p)_{SCL}}{p_p} = \frac{\text{Perimeter under SCL}}{\text{Total perimeter}}$$



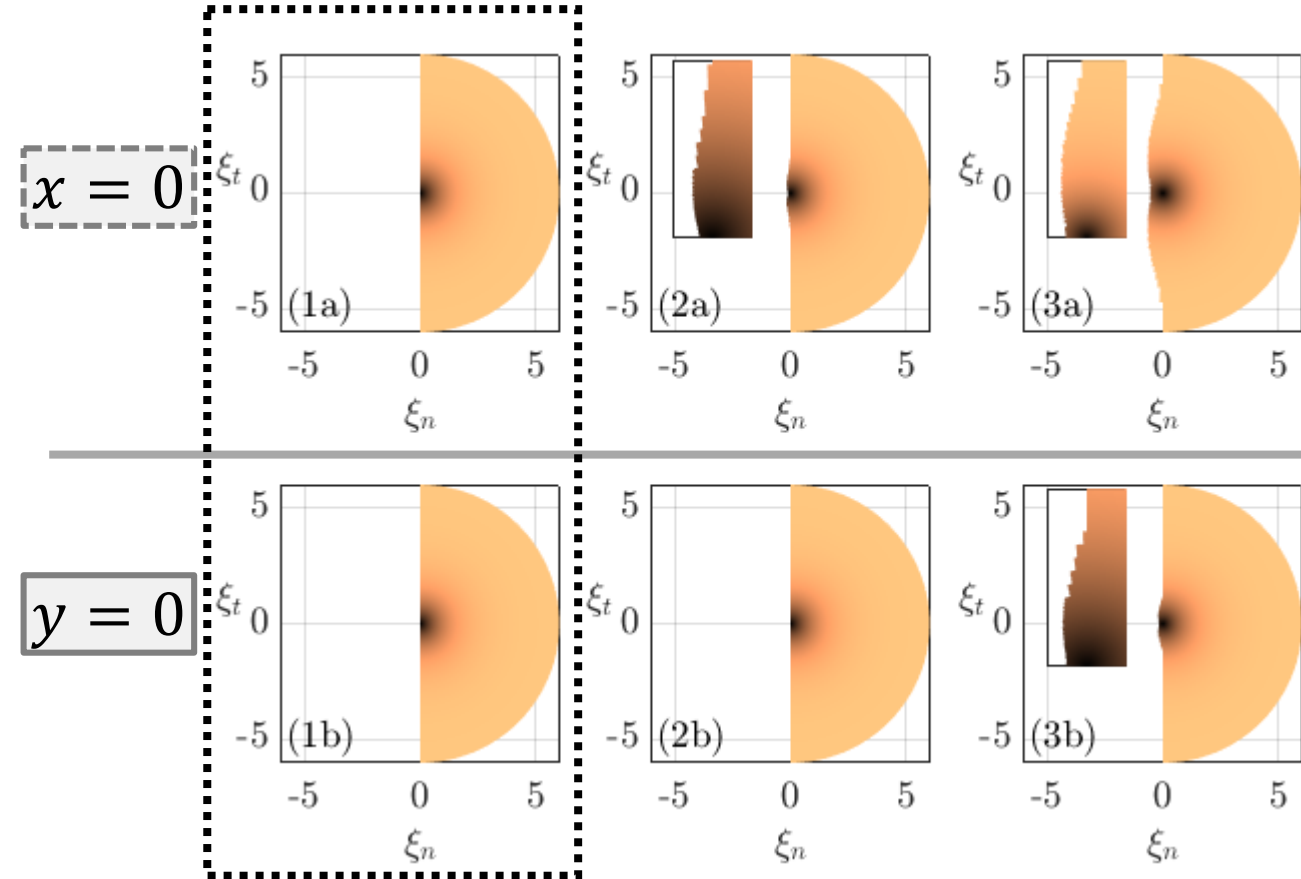
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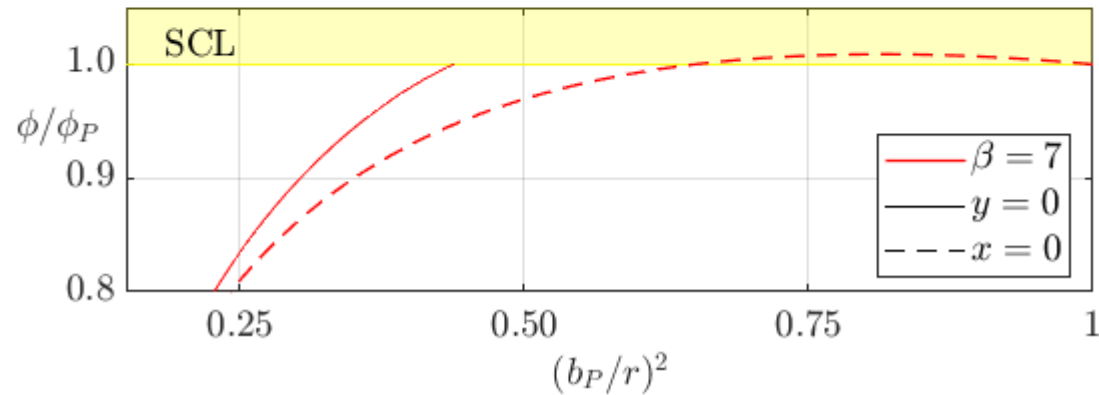
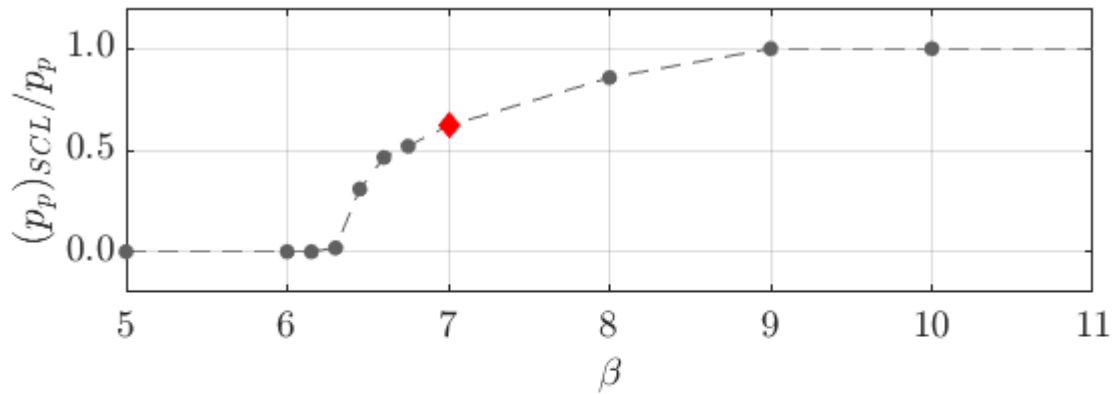
Distribution function of emitted electrons

$\beta = 6$



SPACE CHARGE LIMITED (SCL) TRANSITION

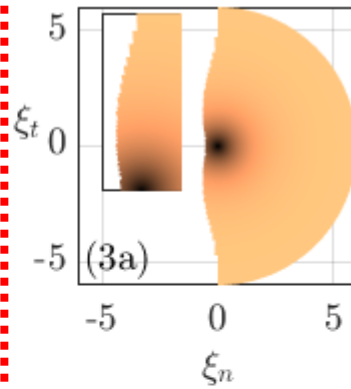
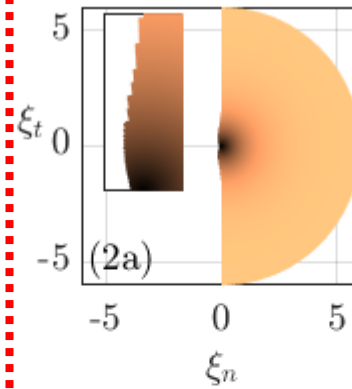
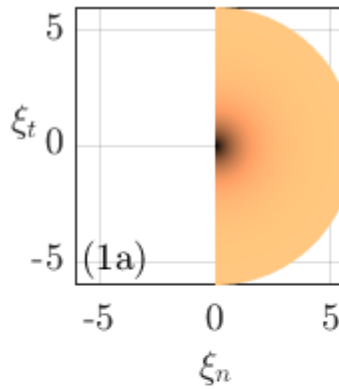
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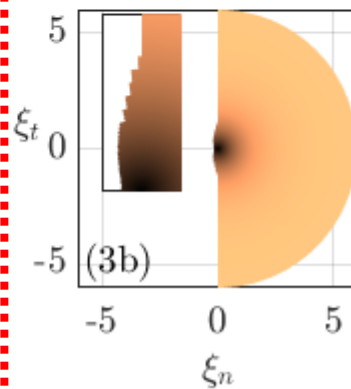
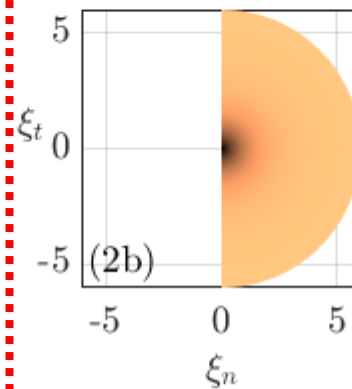
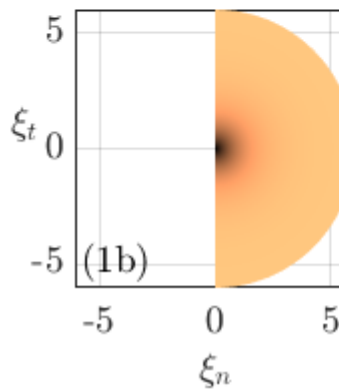
Distribution function of emitted electrons

$\beta = 7$

$x = 0$

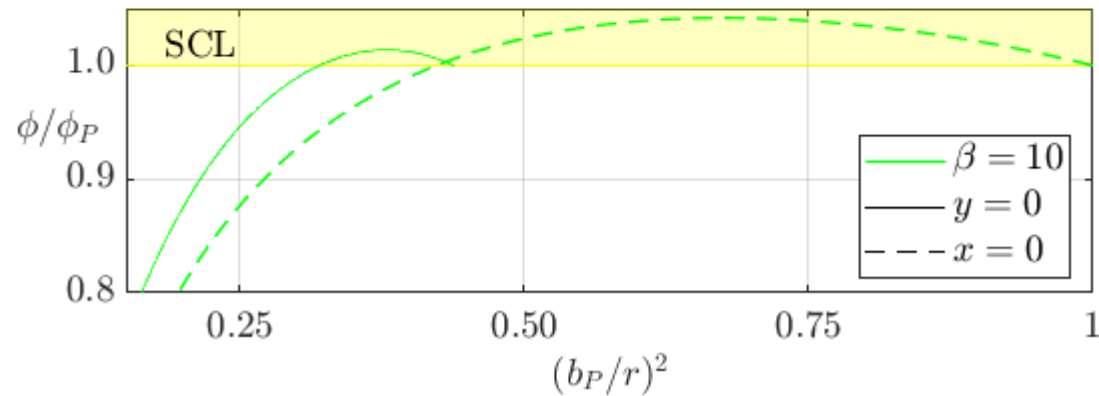
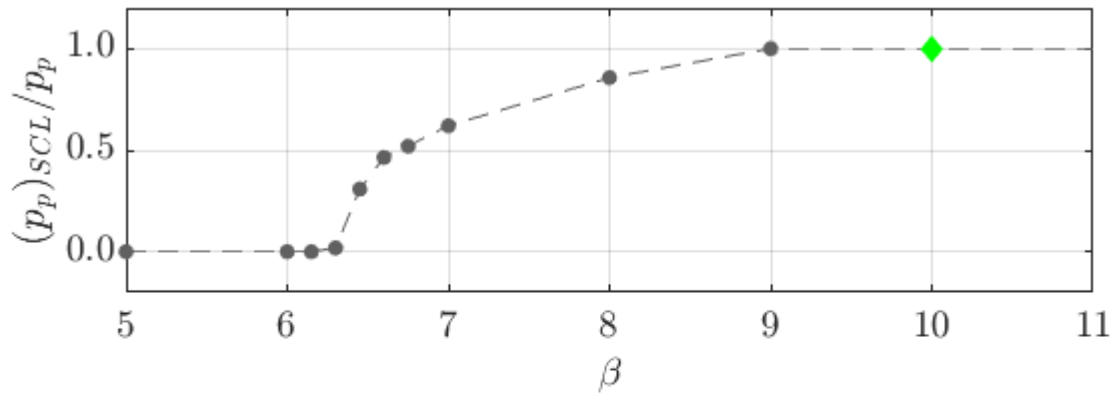


$y = 0$



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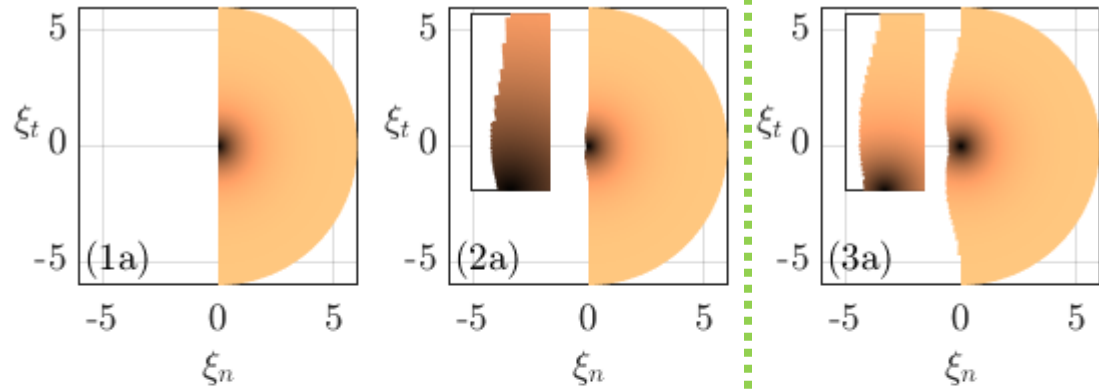
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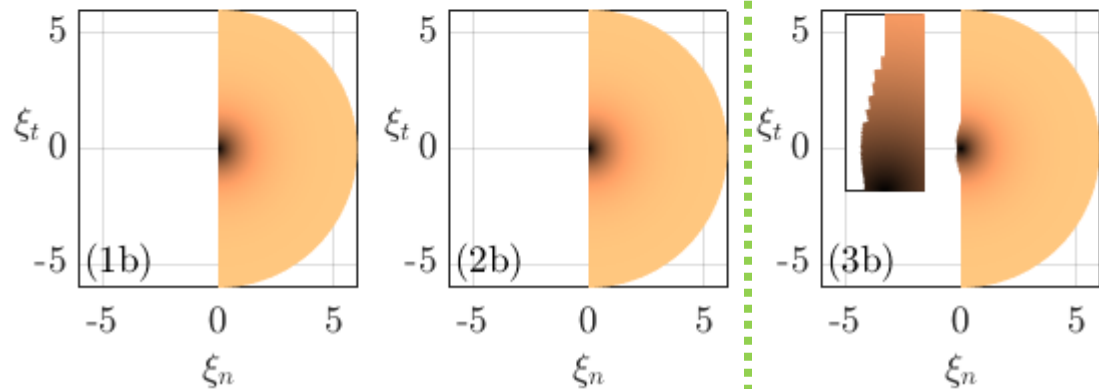
Distribution function of emitted electrons

$\beta = 10$

$x = 0$



$y = 0$



CONCLUSIONS

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- We developed a 2D stationary Vlasov-Poisson solver for the kinetic analysis of emissive objects with elliptical cross-sections
- Numerical and physical limitations of the implementation were discussed (chaotic dynamics, filamentation)
- Transition to SCL regime as a function of the emission level was investigated
- Elliptical probes resemble cylindrical ones with radius equal to the semi-major axis

Find more details at:

- *L. Chiabò, G. Sánchez-Arriaga, “Limitations of Stationary Vlasov-Poisson Solvers in Probe Theory”* (currently under review)
- *L. Chiabò, S. Shahsavani, G. Sánchez-Arriaga, “Kinetic Analysis of Elliptical Emissive Probes Immersed in Maxwellian Plasmas”* (in preparation)



THE END

Thank you for your attention!

L. Chiabò, G. Sánchez-Arriaga

Universidad Carlos III de Madrid, Spain

APS-DPP 2020, 9-13 November